

SIMULTANEOUS AND NON-SIMULTANEOUS BLOW-UP
FOR A NON-LOCAL DIFFUSION SYSTEM

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Abstract: In this work, the non-local diffusion system with Neumann boundary conditions

$$\begin{aligned}u_t(x, t) &= \int_{\Omega} J(x - y)(u(y, t) - u(x, t))dy + f(u(x, t), v(x, t)) \\v_t(x, t) &= \int_{\Omega} J(x - y)(v(y, t) - v(x, t))dy + g(u(x, t), v(x, t)) \\u(x, 0) &= u_0(x), \quad v(x, 0) = v_0(x), \quad x \in \Omega,\end{aligned}\tag{1}$$

is studied, where $(x, t) \in \Omega \times (0, T)$, Ω is a bounded, connected and smooth domain, f and g are continuous functions and $(u_0, v_0) \in C(\overline{\Omega}) \times C(\overline{\Omega})$, nonnegative and real functions. Existence and uniqueness of solutions is proved. For some particular functions f, g , the simultaneous and non-simultaneous blow-up for solutions is analyzed. Finally, the blow-up rate for the solution is given.

AMS Subject Classification: 35K57, 35B40

Key Words: nonlocal diffusion; system equations; Neumann boundary conditions; simultaneous and non-simultaneous blow-up

1. Introduction

The coupled parabolic system

$$u_t = \Delta u + f(u, v), \quad v_t = \Delta v + g(u, v),\tag{2}$$

where $u(x, 0), v(x, 0)$ are nonnegative bounded functions and f, g continuous

Received: February 14, 2023

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functions, has been studied by several authors. The existence and uniqueness of solutions and the blow-up phenomena have been analyzed, see for instance, [8], [12], [13], [15], [16] and [17]. For the case $f(u, v) = u^r v^p$, $g(u, v) = v^q u^s$, whit $p, q, r, s > 0$, Dickstein and Escobedo in [8], studied (2) considering bounded domains. For the case $f(u, v) = u^r + v^p$, $g(u, v) = u^q + v^s$ whit $p, q, r, s > 0$, Souplet and Tayachi in [17], studied the Cauchy problem for (2). Rossi and Souplet in [15], studied (2) considering bounded domains and homogeneous Dirichlet boundary conditions. They show that the non-simultaneous and simultaneous blow-up are possible for the exponent region $r > q + 1$ or $s > p + 1$, both depending on the initial data.

Equations of the form

$$u_t(x, t) = J * u - u(x, t) = \int_{\mathbb{R}^n} J(x - y)u(y, t)dy - u(x, t), \quad (3)$$

and variations of it, have been widely used in the last decade to model diffusion processes, see for instance [1], [2], [6] and [10]. As stated in [10], $J : \mathbb{R}^n \rightarrow \mathbb{R}$ is a non-negative, smooth, symmetric radially and strictly decreasing function, with $\int_{\mathbb{R}^n} J(x)dx = 1$, supported in the unitary ball. If $u(x, t)$ is thought as a density at the point x at time t and $J(x - y)$ is thought as the probability distribution of jumping from location y to location x , then $(J * u)(x, t)$ is the rate at which individuals are arriving to position x from all other places and $-u(x, t) = -\int_{\mathbb{R}^n} J(y - x)u(x, t)dy$ is the rate at which they are leaving location x to travel to all other sites. This consideration, in the absence of external sources, leads immediately to the fact that the density u satisfies equation (3). This equation is called nonlocal diffusion equation because the diffusion of the density u at a point x and time t does not only depend on $u(x, t)$, but also on all the values of u in a neighborhood of x through the convolution term $J * u$. This equation shares many properties with the classical heat equation $u_t = \Delta u$ such as the fact that bounded stationary solutions are constant, a maximum principle holds for both of them and even, if J is compactly supported, perturbations propagate with infinite speed.

Bogoya in [3], studied the problem with Neumann boundary condition and source term

$$\begin{aligned} u_t(x, t) &= \int_{\Omega} J(x - y)(u(y, t) - u(x, t))dy + f(u(x, t)), \quad x \in \Omega, \quad t > 0, \\ u(x, 0) &= u_0(x), \quad x \in \Omega, \end{aligned} \quad (4)$$

where $u_0 \in C(\overline{\Omega})$ is nonnegative function and $\Omega \subset \mathbb{R}^N$ a bounded, smooth and connected domain, f a function of u representing reaction term (source). The

author proves the existence and uniqueness of solutions. For some conditions on f , it is shown that the solution becomes unbounded at time T (blow-up). For blowing-up solutions, the rate of blow-up (that is the speed at which solutions go to infinity an time T) is also analyze. For $f(u) = e^u$, the author considers the radial case in the ball of radius $r > 0$. It is proved that if radially symmetric initial condition has a unique maximum at the origin, then the solution is radially symmetric and has a unique maximum at the origin which it is the only blow up point.

Perez-Llanos and Rossi in [11], studied the problem (4) for $f(u) = u^p$ with $p > 0$, where $u_0 \in C(\overline{\Omega})$ is nonnegative function. They prove that for, nonnegative and nontrivial solutions, blow up in finite time for the solutions, if and only if, $p > 1$. Moreover, they find that the blow-up rate is the same that the one that holds for the ODE $u'(t) = u^p(t)$, that is, $\lim_{t \rightarrow T^-} (T - t)^{1/(p-1)} \|u(\cdot, t)\|_\infty = (1/(p-1))^{1/(p-1)}$.

Bogoya in [4], studied the following nonlocal reaction-diffusion system with Neumann boundary conditions

$$\begin{aligned} u_t(x, t) &= \int_{\Omega} J(x-y)(u(y, t) - u(x, t))dy + v^p(x, t), \quad (x, t) \in \Omega \times (0, T) \\ v_t(x, t) &= \int_{\Omega} J(x-y)(v(y, t) - v(x, t))dy + u^q(x, t), \quad (x, t) \in \Omega \times (0, T) \\ u(x, 0) &= u_0(x), \quad v(x, 0) = v_0(x), \quad x \in \Omega, \end{aligned} \quad (5)$$

with $p, q > 0$, $u_0(x), v_0(x) \in C(\overline{\Omega})$ nonnegative and nontrivial functions and $\Omega \subset \mathbb{R}^N$ ($N \geq 1$) a bounded connected and smooth domain. The author proves the existence and uniqueness of nonnegative solutions (u, v) of (5) and shows that the solution (u, v) is unique if $pq \geq 1$, or if one of the initials conditions is not zero for $pq < 1$. The, the globally existence for the solution of (5) is proved. It is shown that if $pq > 1$ and u_0, v_0 , are nonnegative and nontrivial functions, then the solution (u, v) of (5) blows up in finite time T , if $pq \leq 1$ then the solution (u, v) of (5) exists globally. Finally it is considered the blow-up rate for the solution (u, v) of (5). Bogoya and Gómez in [5] studied the numerical approximation of (5).

General problem: Our first objective in this paper is to study the nonlocal diffusion system

$$\begin{aligned} u_t(x, t) &= \int_{\Omega} J(x-y)(u(y, t) - u(x, t))dy + f(u(x, t), v(x, t)), \\ v_t(x, t) &= \int_{\Omega} J(x-y)(v(y, t) - v(x, t))dy + g(u(x, t), v(x, t)), \\ u(x, 0) &= u_0(x), \quad v(x, 0) = v_0(x), \quad x \in \Omega, \end{aligned} \quad (6)$$

where u_0, v_0 are nonnegative bounded functions, $(x, t) \in \Omega \times (0, T)$, $\Omega \subset \mathbb{R}^N$ is a bounded domain with smooth boundary and $f, g : \mathbb{R}^+ \times \mathbb{R}^+ \rightarrow \mathbb{R}$ are continuous functions. Since the integrating is on Ω , is considered that the diffusion takes place only in Ω , sot that, the individuals may not enter nor leave Ω . This is the analogous of what is called in the literature as homogeneous Neumann boundary conditions.

Sum of power terms: A second objective is to the study of (6) with $f(u, v) = u^r + v^p$, $g(u, v) = u^q + v^s$ for $p, q, r, s > 0$, $u_0, v_0 \in C(\Omega)$, $(x, t) \in \Omega \times (0, T)$,

$$\begin{aligned} u_t(x, t) &= \int_{\Omega} J(x-y)(u(y, t) - u(x, t))dy + u^r(x, t) + v^p(x, t), \\ v_t(x, t) &= \int_{\Omega} J(x-y)(v(y, t) - v(x, t))dy + u^q(x, t) + v^s(x, t), \\ u(x, 0) &= u_0(x), \quad v(x, 0) = v_0(x), \quad x \in \Omega. \end{aligned} \quad (7)$$

The system (7), can be viewed as a combination of the following two systems: the problem (5) and

$$\begin{aligned} u_t(x, t) &= \int_{\Omega} J(x-y)(u(y, t) - u(x, t))dy + u^r(x, t), \quad (x, t) \in \Omega \times (0, T) \\ v_t(x, t) &= \int_{\Omega} J(x-y)(v(y, t) - v(x, t))dy + v^s(x, t), \quad (x, t) \in \Omega \times (0, T) \\ u(x, 0) &= u_0(x), \quad v(x, 0) = v_0(x), \quad x \in \Omega, \end{aligned} \quad (8)$$

(8) is uncoupled system and non-simultaneous blowing-up solutions exists, see [11]. In [3], it is studied the coupled system (5), which has only simultaneous blowing-up solutions if $pq > 1$. Therefore it is natural to ask whether the blow-up is simultaneous or not for (7).

It is said that a solution (u, v) blows up in finite time if only if there exists a finite time $T > 0$, such that

$$\limsup_{t \nearrow T} (\|u(x, t)\|_{L^\infty(\Omega)} + \|v(x, t)\|_{L^\infty(\Omega)}) = \infty.$$

If $T = \infty$, the solution (u, v) is global, i.e. the solution exists for all $t \geq 0$.

It is said that the blow-up is simultaneous if

$$\limsup_{t \nearrow T} \|u(x, t)\|_{L^\infty(\Omega)} = \limsup_{t \nearrow T} \|v(x, t)\|_{L^\infty(\Omega)} = \infty.$$

One can note a priori that there is not a reason for both components of the system to blow up simultaneously. In fact, it could happen that one of the

components blows up as $t \rightarrow T$, while the other, remains bounded on $[0, T)$. This phenomenon is called non-simultaneous blow up.

The following information is necessary for the study: If (α, β) is the unique solution of

$$\begin{pmatrix} r-1 & p \\ q & s-1 \end{pmatrix} \begin{pmatrix} \alpha \\ \beta \end{pmatrix} = \begin{pmatrix} -1 \\ -1 \end{pmatrix}, \quad (9)$$

then $\alpha = (1 - s + p)/\Delta$, $\beta = (1 - r + q)/\Delta$, $\Delta = (r - 1)(s - 1) - pq \neq 0$.

The rest of the paper is organized as follows: In Section 2, we study the general problem (6), the existence and uniqueness of nonnegative solutions (u, v) as well as a comparison principle for the solutions is studied. In Section 3, we study the sum of power terms (7), the global existence and blow-up of solutions is proved, non-simultaneous blows up and the blow-up rate is studied.

2. General problem

In this section, we prove the existence and uniqueness of nonnegative solutions (u, v) of (6) by following the ideas of [8].

Let $t_0 > 0$ be fixed and

$$X_{t_0} = C([0, t_0] : C(\overline{\Omega}) \times C(\overline{\Omega})) = \{(u, v) : [0, t_0] \rightarrow C(\overline{\Omega}) \times C(\overline{\Omega}) : (u, v) \text{ is continuous}\},$$

a Banach space with the norm

$$\begin{aligned} |||(u, v)||| &= \max_{0 \leq t \leq t_0} \|(u(\cdot, t), v(\cdot, t))\|_I, \text{ with } I = L^\infty(\overline{\Omega}) \times L^\infty(\overline{\Omega}), \\ \|(u(\cdot, t), v(\cdot, t))\|_I &= \max_{x \in \overline{\Omega}} |u(x, t)| + \max_{x \in \overline{\Omega}} |v(x, t)|. \end{aligned}$$

Let $\|u\|_\kappa = \max_{0 \leq t \leq t_0} \|u(\cdot, t)\|_{L^\infty(\overline{\Omega})}$.

Let $P_{t_0} = \{(u, v) \in X_{t_0} : u \geq 0, v \geq 0\}$ be a closed subspace of X_{t_0} . We define the operator $\psi : P_{t_0} \rightarrow P_{t_0}$ as $\psi_{(u_0, v_0)}(u, v) = (T_{u_0}(u), S_{v_0}(v))$, where

$$\begin{aligned} T_{u_0}(u)(x, t) &= \int_0^t \int_\Omega J(x - y)(u(y, s) - u(x, s)) dy ds \\ &\quad + \int_0^t f(u(x, s), v(x, s)) ds + u_0(x), \\ S_{v_0}(v)(x, t) &= \int_0^t \int_\Omega J(x - y)(v(y, s) - v(x, s)) dy ds \\ &\quad + \int_0^t g(u(x, s), v(x, s)) ds + v_0(x). \end{aligned} \quad (10)$$

We begin the study considering the functions f and g locally Lipschitz. The following lemma is very important for the study and its proof is analogous to that given in [4], reason why, we omit here.

Lemma 1. *Let f and g be locally Lipschitz function, $(u_0, v_0), (w_0, z_0) \in C(\overline{\Omega}) \times C(\overline{\Omega})$ and $(u, v), (w, z) \in P_{t_0}$. Then there exists a positive constant $C = C(K_1, K_2, \Omega, J)$ such that*

$$\begin{aligned} |||\psi_{(u_0, v_0)}(u, v) - \psi_{(w_0, z_0)}(w, z)||| &\leq Ct_0 |||(u, v) - (w, z)||| \\ &+ \|(u_0, v_0) - (w_0, z_0)\|_I. \end{aligned} \quad (11)$$

Theorem 2. *Let f and g be locally Lipschitz function and $(u_0, v_0) \in C(\overline{\Omega}) \times C(\overline{\Omega})$, nonnegative and real functions, then, there exists a unique solution (u, v) of (6) such that $(u, v) \in P_{t_0}$. Moreover, (u, v) can be extended to a maximal interval $[0, T)$ with $T \leq \infty$.*

Proof. Following ideas of the proof of Lemma 1, we see that the operator $\psi : P_{t_0} \cap B_r(0, 0) \rightarrow P_{t_0} \cap B_r(0, 0)$ is well defined, where

$$r = \max\{\|u_0\|, \|v_0\|\} + 1 \quad \text{and} \quad B_r(0, 0) = \{(u, v) : \|(u, v)\|_I < r\}. \quad (12)$$

Now, taking $(u_0, v_0) = (w_0, z_0)$ in Lemma 1 and choosing t_0 such that $Ct_0 < 1$, we obtain that $\psi(u, v)$ is a strict contraction of $P_{t_0} \cap B_r(0, 0)$ into itself, therefore, there exists a unique fixed point (u, v) of $\psi(u, v)$ in $P_{t_0} \cap B_r(0, 0)$ by the Banach fixed point theorem, which is the uniqueness of solution to (6), in $\overline{\Omega} \times [0, t_0]$. Uniqueness implies that the solution can be extended to a maximal interval $[0, T)$, with $T \leq \infty$.

Remark 3. The solution (u, v) of (6) depend continuously on the initial data. In fact if (u, v) and (w, z) are solutions to (6) with initial data (u_0, v_0) and (w_0, z_0) respectively, then there exists a constant $\tilde{C} = \tilde{C}(t_0, K_3, \Omega, J)$ such that

$$|||(u, v) - (w, z)||| \leq \tilde{C} \|(u_0, v_0) - (w_0, z_0)\|_I.$$

Remark 4. $(u, v) \in P_{t_0}$ is a solutions of (6),

$$\begin{aligned} u(x, t) &= \int_0^t \int_{\Omega} J(x-y)(u(y, s) - u(x, s)) dy ds \\ &\quad + \int_0^t f(u(x, s), v(x, s)) ds + u_0(x), \\ v(x, t) &= \int_0^t \int_{\Omega} J(x-y)(v(y, s) - v(x, s)) dy ds \\ &\quad + \int_0^t g(u(x, s), v(x, s)) ds + v_0(x). \end{aligned} \tag{13}$$

Following the ideas of [8], the existence and uniqueness of the solution of (6) is studied under conditio that f and g , are continuous functions.

Theorem 5. *Let f and g continuous functions and $(u_0, v_0) \in C(\overline{\Omega}) \times C(\overline{\Omega})$ nonnegative and real functions, then, there exists a unique solution (u, v) of (6), such that $(u, v) \in P_{t_0}$.*

We will use the notation $(a, b) \geq (c, d)$ to indicate that $a \geq c$ and $b \geq d$.

Definition 6. Let $\overline{u}, \overline{v} \in C^1([0, T]; C(\overline{\Omega}))$. $(\overline{u}, \overline{v})$ is called a supersolution of (6) if

$$\begin{aligned} \overline{u}_t(x, t) &\geq \int_{\Omega} J(x-y)(\overline{u}(y, t) - \overline{u}(x, t)) dy + f(\overline{u}(x, t), \overline{v}(x, t)) \\ \overline{v}_t(x, t) &\geq \int_{\Omega} J(x-y)(\overline{v}(y, t) - \overline{v}(x, t)) dy + g(\overline{u}(x, t), \overline{v}(x, t)) \\ \overline{u}(x, 0) &\geq u_0(x), \quad \overline{v}(x, 0) \geq v_0(x), \quad x \in \Omega. \end{aligned} \tag{14}$$

Analogously, $(\underline{u}, \underline{v})$, is called a subsolution of (6), if it satisfies the opposite inequalities.

Now, we consider the following hypothesis on f and g :

H_1 : f and g are increasing functions in each variable, i.e. $f(u_2, v_2) \geq f(u_1, v_1)$ and $g(u_2, v_2) \geq g(u_1, v_1)$ for all $0 \leq u_1 \leq u_2$, $0 \leq v_1 \leq v_2$.

The hypothesis H_1 allows us to establish some comparison results, which are given in the following lemma. Its proof is analogous to that given in Lemma 2.2 of [8], a reason why we omit it here.

Lemma 7. *Let f, g , be continuous functions satisfying H_1 , and let (\bar{u}, \bar{v}) , a supersolution of (6) in $\Omega \times (0, T)$. Then, there exists a solution (u, v) of (6) in $\Omega \times (0, T)$ such that $(u, v) \leq (\bar{u}, \bar{v})$. Moreover, if $(\underline{u}, \underline{v})$ is a subsolution of (6) in $\Omega \times (0, T)$, verifying $(\underline{u}, \underline{v}) \leq (\bar{u}, \bar{v})$, then, there exists a solution (u, v) defined on $\Omega \times (0, T)$ such that $\underline{u} \leq u \leq \bar{u}$, $\underline{v} \leq v \leq \bar{v}$.*

As a consequence of these previous results, we have the following corollary, which proof is omitted.

Corollary 8. *Let f and g be locally Lipschitz functions satisfying H_1 and (u, v) the solution of (6), in $\Omega \times (0, T)$.*

1. *If (\bar{u}, \bar{v}) is a supersolution of (6), then $(u, v) \leq (\bar{u}, \bar{v})$ in $\Omega \times (0, T)$.*
2. *If $(\underline{u}, \underline{v})$ is a subsolution of (6), then $(\underline{u}, \underline{v}) \leq (u, v)$ in $\Omega \times (0, T)$.*

Definition 9. A solution (u_M, v_M) of (6) in $\Omega \times (0, T)$ is a maximal solution of (6), if given any other solution (u, v) of (6), we have $(u, v) \leq (u_M, v_M)$ in $\Omega \times (0, T)$.

Lemma 10. *Let f, g be continuous functions satisfying H_1 . Then, (6) has a maximal solution (u, v) in $\Omega \times (0, T)$.*

Proof. As in Theorem 5, let $(f_n)_n$ and (g_n) decreasing sequences of locally Lipschitz functions, with $f_n \rightarrow f$, $g_n \rightarrow g$ as $n \rightarrow \infty$. Moreover, we can take f_n, g_n satisfying the condition H_1 , for all $n \in N$. Let (u_n, v_n) the solution of (1), then $u_M(x, t) = \lim_{n \rightarrow \infty} u_n(x, t)$, $v_M(x, t) = \lim_{n \rightarrow \infty} v_n(x, t)$, and by Remark (4), we obtain what (u_M, v_M) , is a solution of (6), in $\Omega \times (0, T)$. We claim that (u_M, v_M) is the maximal solution of (6). Indeed, if (u, v) is any other solution of (6) with source term f, g then (u, v) is a subsolution of (6) with source term f_n, g_n for all $n \in N$ as $(f_n)_n$ and (g_n) decreasing sequences, therefore $(u, v) \leq (u_n, v_n)$ for all $n \in N$. Letting $n \rightarrow \infty$, we get $(u, v) \leq (u_M, v_M)$ in $(x, t) \in \Omega \times (0, T)$.

Remark 11. If (\bar{u}, \bar{v}) is a positive supersolution of (6) and $(\underline{u}, \underline{v})$ is a subsolution of (6), then $(\bar{u}, \bar{v}) \geq (\underline{u}, \underline{v})$.

The following lemma is very important for the case when f and g are non-Lipchitz functions. The demonstration is very similar to this given in [8], so that we omit here.

Lemma 12. *Let f, g be continuous functions satisfying H_1 , with $f(0, 0) = g(0, 0) = 0$ and $f \neq 0, g \neq 0$. Let $(u_0, v_0) = (0, 0)$ in Ω and (\bar{u}, \bar{v}) a supersolution of (6), such that $(\bar{u}, \bar{v}) > (0, 0)$ in $\Omega \times (0, T)$. Then $(\bar{u}, \bar{v}) > (u, v)$ in $\Omega \times (0, T)$ for any solution (u, v) of (6), with $(u_0, v_0) = (0, 0)$.*

3. General problem

Now, we analyze the problem (7). As $p, q, r, s > 0$, we note that $f(u, v) = u^r + v^p$, $g(u, v) = u^q + v^s$ satisfy H_1 .

3.1 Existence and uniqueness

Theorem 13. *Let $u_0, v_0 \in C(\Omega)$ with $(u_0, v_0) \neq (0, 0)$ in Ω . Then there exists a unique solution (u, v) of (7) such that $(u, v) \in P_{t_0}$.*

Proof. It is obtained from Theorem 5.

Theorem 14. *Let $u_0, v_0 \in C(\Omega)$ with $(u_0, v_0) = (0, 0)$ for all $x \in \Omega$. If $r \geq 1, s \geq 1$ and $pq \geq 1$, then the problem (7), has only the trivial solution. Otherwise, there exists a unique positive solution (u, v) of (7) in $(x, t) \in \Omega \times (0, T)$, which is the maximal.*

Remark 15. The special case $r = s = 0$ leads to study the problem (5), which is studied in [4] and a similar result to Theorem 13 is obtained.

Remark 16. By a previous remark, it follows from the proof of Theorem 14 that if $r \geq 1, s \geq 1$, and $pq \geq 1$, then, there is no nontrivial subsolution leaving $(0, 0)$.

3.2 Blow-up analysis

In this section, we study conditions under which, the solution of (7) exists globally or blow-up. In what follows, we build on ideas developed in [17].

Theorem 17. *If $\max\{r, s, pq\} > 1$, the solutions of (7), blow-up in finite time.*

Proof. We shall proceed in two cases.

Case 1. Let $pq = \max\{r, s, pq\} > 1$. Let $(\underline{u}, \underline{v})$ the solution of (5), which blow-up in finite time $T > 0$ for $pq > 1$, see [4]. As $(\underline{u}, \underline{v})$ is nonnegative then is a subsolution of (7) in $\Omega \times [0, T)$. If $(0, 0) \neq (\underline{u}_0(x), \underline{v}_0(x)) \leq (u_0(x), v_0(x))$ in Ω , then by Lemma 7 we have that the solution (u, v) of (7) blow-up in finite time.

Case 2. Let $r = \max\{r, s, pq\} > 1$ or $s = \max\{r, s, pq\} > 1$, these two cases can be handled with the same argument, so that, we consider $r > 1$.

Let \underline{u} be the nonnegative solution of (4) with $f(u) = u^r$. As $r > 1$, we have that \underline{u} blow-up in finite time, see [11]. Let \underline{v} , the nonnegative solution of (4) with $f = 0$, see [7]. Moreover, we have that $(\underline{u}, \underline{v})$ is a subsolution of (7) in $\Omega \times [0, T)$. If $(0, 0) \neq (\underline{u}_0(x), \underline{v}_0(x)) \leq (u_0(x), v_0(x))$ in Ω , then by Lemma 7 we have that the solution (u, v) of (7) blow-up in finite time.

Theorem 18. *If $\max\{r, s, pq\} \leq 1$, then all solutions of (7) are global.*

Proof. Let $\max\{r, s, pq\} \leq 1$ and α, β, Δ , given in (9). Let $\alpha_1 = 2\alpha$, $\beta_1 = 2\beta$. We have $\alpha_1, \beta_1 > 0$, $\alpha_1(1-r) - \beta_1 p = 2$ and $\beta_1(1-s) - \alpha_1 q = 2$. Let $\Delta > 0$ and $C > 0$, we consider the functions $\bar{u} = (t+C)^{\alpha_1}$, $\bar{v} = (t+C)^{\beta_1}$ in $\Omega \times [0, \infty)$. We have

$$\begin{aligned} \bar{u}_t - \int_{\Omega} J(x-y)(\bar{u}(y, t) - \bar{u}(x, t))dy - \bar{u}^r - \bar{v}^p \\ = \alpha_1(t+C)^{\alpha_1-1} - (t+C)^{\alpha_1-\beta_1 r-2} - (t+C)^{\alpha_1-\alpha_1 q-2} \geq 0, \end{aligned} \quad (15)$$

for C sufficiently large. Similarly, we have

$$\bar{v}_t - \int_{\Omega} J(x-y)(\bar{v}(y, t) - \bar{v}(x, t))dy - \bar{u}^q - \bar{v}^s \geq 0.$$

Therefore (\bar{u}, \bar{v}) , is a supersolution of (7). Let $(u_0, v_0) \geq (0, 0)$ with $(u_0, v_0) \leq (\bar{u}(x, 0), \bar{v}(x, 0))$. By Lemma 7 and the uniqueness the result of the Theorems 13 and 14, we have the solution (u, v) of (7) is global.

If $\Delta = 0$, then $(1-r)(1-s) = pq$. Let $\bar{u} = Ce^{\mu t}$, $\bar{v} = Ce^{\nu t}$ in $\Omega \times [0, \infty)$ such that $\nu = (1-r)p^{-1}\mu = q(1-s)^{-1}\mu$. We have

$$\begin{aligned} \bar{u}_t - \int_{\Omega} J(x-y)(\bar{u}(y, t) - \bar{u}(x, t))dy - \bar{u}^r - \bar{v}^p = \mu\bar{u} - \bar{u}^r - \bar{v}^p \\ = e^{\mu t}(\mu C - C^r e^{-\nu p t} - C^p e^{-\mu r t}) \geq e^{\mu t}(\mu C - C^r - C^p) \geq 0, \end{aligned} \quad (16)$$

for μ sufficiently large. Similarly, we have

$$\bar{v}_t - \int_{\Omega} J(x-y)(\bar{v}(y, t) - \bar{v}(x, t))dy - \bar{u}^q - \bar{v}^s \geq e^{\nu t}(\nu C - C^q - C^s) \geq 0. \quad (17)$$

Therefore (\bar{u}, \bar{v}) is a supersolution of (7) and the same conclusion for the case $\Delta > 0$ holds.

Next, we will study the conditions under which non-simultaneous blow-up the solution of (7) could occur. Consider the system of ODE

$$\begin{aligned} u'(t) &= u^r(t) + v^p(t), \quad v'(t) = u^q(t) + v^s(t), \quad t > 0 \\ u(0) &= a, \quad v(0) = b. \end{aligned} \quad (18)$$

Remark 19. 1. If $p = q = r = s = 1$, then $u(t) = C_1 e^{2t}$, $v(t) = C_2 e^{2t}$ is solution of (18).

2. Let $pq > 1$ and (u, v) be nonnegative solutions of (18). We have $u'(t) \geq v^p(t)$, $v'(t) \geq u^q(t)$, therefore (u, v) is a supersolution of $w'(t) = z^p(t)$, $z'(t) = w^q(t)$, $w(0) = c \leq a$, $z(0) = d \leq b$. As the solution (w, z) , simultaneous blow-up in finite time $T > 0$ for $qr > 1$, then by Comparison Principle, (u, v) simultaneous blow-up in finite time $T > 0$.

Remark 20. $(u(t), v(t))$ is a flat solution (a solution that does not depend on x) of (7), with initial datum $u(x, 0) = a \geq 0$, $v(x, 0) = b \geq 0$ if only if $(u(t), v(t))$ is a solution of the system of ODE (18).

The following proposition refers to the flat solution (7) and is analogous to Proposition 2.2(i) of [17], reason for which we omit its demonstration.

Proposition 21. Let $r > q + 1$. If $c_1 u_0^{1+q-r} + v_0 < c_2 u_0^{(r-1)/(s-1)}$, with $c_1 = \frac{1}{r-1-q}$, $c_2 = \left(\frac{r-1}{s-1}\right)^{1/(s-1)}$, then, for the flat solution $(u(t), v(t))$ of (7), must be u blow-up in finite time $T > 0$, while v remains bounded on $[0, T)$, i.e. satisfies

$$\lim_{t \nearrow T} u(t) = \infty \quad \text{and} \quad \sup_{t \in (0, T)} v(t) < \infty. \quad (19)$$

For $s > p + 1$, the analogue obviously holds by exchanging the roles of u, v .

Theorem 22. (Non-simultaneous Blow-up) Let (u, v) be a positive blowing-up solution of (7). If $r > q + 1$ or $s > p + 1$, then there exist u_0 and v_0 such that non-simultaneous blow-up occurs.

Proof. Assume that $r > q + 1$, then by Proposition 21, there exists u_0 and v_0 such that $\lim_{t \nearrow T} \|u(t)\|_\infty = \infty$ and $\sup_{\Omega \times (0, T)} v(x, t) < \infty$ for some finite time

$T > 0$. Hence, non-simultaneous blow-up occurs. For $s > p + 1$, the analogue of Proposition 21 holds by exchanging the roles of u and v .

3.3 Blow-up rates

In this section, we analyze the blow-up rate of the solutions of (7). We assume that $x = 0 \in \Omega$, and note that for smooth radially symmetric and nondecreasing initial conditions (that is, when $u_0(r)$, $v_0(r)$ are C^1 such that $u'_0(r) \leq 0$, $v'_0(r) \leq 0$) the solutions are also radially symmetric and radially nondecreasing (that is, it holds that $u_r(r, t) \leq 0$, $v_r(r, t) \leq 0$). Hence, for every $t \in (0, T)$, the maximum of both components is attained at $x = 0$. We state this result as follows. For a proof, we refer to Lemma 4.1 in [11].

Lemma 23. *If $\Omega = B(0; R)$ is a ball and (u_0, v_0) are smooth, radially symmetric, and nondecreasing initial conditions (i.e. $u_0(r)$, $v_0(r)$ are C^1 such that $u'_0(r) \leq 0$, $v'_0(r) \leq 0$) then both components u , v of the solution of (7) are radially symmetric and radially nondecreasing (they verify $u_r(r, t) \leq 0$, $v_r(r, t) \leq 0$ for every $r \in [0, R]$ and every $t > 0$).*

Next, we study the blow-up rate of the solutions of (7).

Theorem 24. *Let (u, v) be a positive blowing-up solution of (7), such that the maximum is reached at $x = 0$ for all $t \in (0, T)$. Let $r < \frac{p(q+1)}{(p+1)}$ and $s < \frac{q(p+1)}{(q+1)}$, then there exists C_1 , C_2 , C_3 , C_4 positive constants such that*

$$\begin{aligned} C_1(T-t)^{-\mu} &\leq u(0, t) \leq C_2(T-t)^{-\mu}, \quad 0 < t < T, \\ C_3(T-t)^{-\nu} &\leq v(0, t) \leq C_4(T-t)^{-\nu}, \quad 0 < t < T, \end{aligned} \tag{20}$$

with $\mu = \frac{p+1}{pq-1}$, $\nu = \frac{q+1}{pq-1}$.

Proof. Let $u(0, t) = \max_{x \in \overline{\Omega}} u(x, t)$ and $v(0, t) = \max_{x \in \overline{\Omega}} v(x, t)$. By (7), we have

$$\begin{aligned} u_t(0, t) &= \int_{\Omega} J(0-y)(u(y, t) - u(0, t))dy + u^r(0, t) + v^p(0, t) \\ &\leq u^r(0, t) + v^p(0, t), \\ v_t(0, t) &= \int_{\Omega} J(0-y)(v(y, t) - v(0, t))dy + u^q(0, t) + v^s(0, t) \\ &\leq u^q(0, t) + v^s(0, t). \end{aligned} \tag{21}$$

As $1 = \int_{R^N} J(\zeta) d\zeta \geq \int_{\Omega} J(\zeta) d\zeta$ and (u, v) is a positive solution, we have

$$u_t(0, t) \geq -u(0, t) + v^p(0, t), \quad v_t(0, t) \geq -v(0, t) + u^q(0, t).$$

Therefore, we have that for all $0 < t < T$

$$-u(0, t) + v^p(0, t) \leq u_t(0, t) \leq u^r(0, t) + v^p(0, t) \quad (22)$$

and

$$-v(0, t) + u^q(0, t) \leq v_t(0, t) \leq u^q(0, t) + v^s(0, t). \quad (23)$$

Multiplying the second inequality of (22) by $u^q(0, t)$ and the first inequality of (23) by $v^p(0, t)$, we have

$$u_t(0, t)u^q(0, t) \leq u^{r+q}(0, t) + v_t(0, t)v^p(0, t) + v^{p+1}(0, t),$$

which is equivalent to

$$\left(\frac{u^{q+1}(0, t)}{q+1} \right)_t - u^{r+q}(0, t) \leq \left(\frac{v^{p+1}(0, t)}{p+1} \right)_t + v^{p+1}(0, t).$$

Multiplying the inequality by $(p+1)e^{(p+1)t}$ and integrating on $[0, t]$ with $t < T$, we have

$$u^q(0, t) \leq C(v(0, t))^{(p+1)q/(q+1)}. \quad (24)$$

Replacing the second inequality of (23) by the inequality (24) and as $s < \frac{q(p+1)}{(q+1)}$, we have

$$\begin{aligned} v_t(0, t) &\leq C(v(0, t))^{(p+1)q/(q+1)} + v^s(0, t) \\ &\leq C(v(0, t))^{(p+1)q/(q+1)} + (v(0, t))^{(p+1)q/(q+1)}. \end{aligned}$$

Therefore,

$$v_t(0, t) \leq (C+1)(v(0, t))^{(p+1)q/(q+1)}.$$

Integrating the inequality from above on $[t, T]$, we obtain that $v(0, t) \geq C_3(T-t)^{-\nu}$, where $\nu = \frac{q+1}{pq-1}$. In analogous way, we obtain $u(0, t) \geq C_1(T-t)^{-\mu}$, where $\mu = \frac{p+1}{pq-1}$.

Doing a similar analysis to the one developed above, we obtain that there exists a constant $C > 0$, such that, for $0 < t < T$

$$C(v(0, t))^{(p+1)q/(q+1)} \leq u^q(0, t). \quad (25)$$

Replacing the first inequality of (23), by the inequality (25) and as $pq > 1$ we have $(p+1)q/(q+1) > 1$ and

$$C(v(0, t))^{(p+1)q/(q+1)} \leq -v(0, t) + C(v(0, t))^{(p+1)q/(q+1)} \leq v_t(0, t).$$

Integrating the inequality from above on $[t, T)$, we obtain

$$v(0, t) \leq C_4(T - t)^{-\nu}.$$

In analogous way, we obtain

$$u(0, t) \leq C_2(T - t)^{-\mu}.$$

Theorem 25. *Let $r > q + 1$ or $s > p + 1$ and (u, v) be a positive solution of (7).*

(i) *If u blows up, then there exists C_1, C_2 positive constants, such that*

$$C_1(T - t)^{-1/(r-1)} \leq \|u(t)\|_\infty \leq C_2(T - t)^{-1/(r-1)}, \quad 0 < t < T.$$

(ii) *If v blows up, then there exists C_3, C_4 positive constants such that*

$$C_3(T - t)^{-1/(s-1)} \leq \|v(t)\|_\infty \leq C_4(T - t)^{-1/(s-1)}, \quad 0 < t < T.$$

Proof. (i) Let $r > q + 1$, then by Theorem 22 we have that non-simultaneous blow-up occurs. Let u blow up in finite time $T > 0$ with $\lim_{t \rightarrow T^-} \|u(t)\|_\infty = \infty$ and v remains bounded, there exists a constant $C > 0$, such that $v(x, t) \leq C$ for all $(x, t) \in \Omega \times (0, \infty)$. Therefore u a nonnegative solution of the problem

$$\begin{aligned} u_t(x, t) &= \int_{\Omega} J(x - y)(u(y, t) - u(x, t))dy + u^r + b(x, t), \\ &\quad (x, t) \in \Omega \times (0, T), \end{aligned} \tag{26}$$

$$u(x, 0) = u_0, \quad x \in \Omega,$$

where $b(x, y) \leq C$, is a positive function. As $r > 1$ and using a similar argument given in [11], we have there exists C_2 a positive constant, such that $\|u(t)\|_\infty \leq C_2(T - t)^{-1/(r-1)}$, $0 < t < T$.

Let $U(t) = \sup_{x \in \Omega} u(x, t)$, $V(t) = \sup_{x \in \Omega} v(x, t)$, by Remark 4 we have for $0 < t < t_1 < t$

$$U(t_1) \leq U(t) + \int_t^{t_1} U^r(s)ds + \int_t^t V^p(s)ds \leq U(t) + \int_t^{t_1} U^r(s)ds + C.$$

As u blows up, there exists a first $t_1 \in [t, T)$, such that $U(t_1) = 2U(t)$. It follows that $U(t_1) = 2U(t) \leq U(t) + (t_1 - t)(2U(t))^r + C$. Therefore $U(t) \leq (T - t)(2U(t))^r + C$ for t close to T , we have $U(t) \geq 2C$. Then $U(t) \geq \frac{1}{2^r}(T - t)^{-1/(r-1)}$. Therefore, $C_1(T - t)^{-1/(r-1)} \leq \|u(t)\|_\infty$. The proof of (ii) proceeds in an analogous way.

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