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POSITIVE SOLUTIONS FOR A SECOND ORDER EXTENDED FISHER-KOLMOGOROV'S EQUATION

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Abstract: We consider the existence of positive solutions of extended Fisher-Kolmogorov second order differential equation. Using a variational method and an approach of Verzini, we obtain the positive bounded solutions of this ODE.

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1. Introduction

In the present paper we study the existence of positive solutions of a secondorder ordinary differential equation (ODE)

$$u'' + cu' + f(t, u) = 0, \quad t \in (a, +\infty),$$
 (1)

coupled with the boundary conditions

$$u(a) = u(+\infty) = 0, (2)$$

where c>0 is a constant, $a\in\mathbb{R}$. We suppose that $f(t,s):\mathbb{R}^2\to\mathbb{R}$ and

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 $f_s(t,s) = \frac{\partial f}{\partial s}(t,s), f_t(t,s) = \frac{\partial f}{\partial t}(t,s)$ are continuous functions which satisfy the following conditions:

$$|c_1|s|^{1+q} \le |f(t,s)| \le |c_2|s|^{1+q}, sf(t,s) \ge 0, \quad \forall (t,s) \in \mathbb{R}^2,$$
 (H1)

$$f_s(t,s)s^2 - Af(t,s) \ge 0, \quad \forall (t,s) \in \mathbb{R}^2,$$
 (H2)

$$f(t,s)s - (2+\alpha)F(t,s) \ge 0, \quad \forall (t,s) \in \mathbb{R}^2, \tag{H3}$$

$$c_3|s|^{q+2} \ge f_t(t,s) \ s \ge 0, \quad \forall (t,s) \in \mathbb{R}^2,$$
 (H4)

where A > 1, $\alpha > 0$, $c_j > 0$, j = 1, 2, 3, and $F(t, s) := \int_{0}^{s} f(t, \tau) d\tau$.

An example of a function f, which satisfies these conditions is $f(t,s) = C|s|^q$.s with $A \in (1, q+1]$, $\alpha = q$, $c_1 = c_2 = c_3 = C > 0$.

We will look for positive solutions u of (1.1) such that $u \in H_0^1(a, +\infty) \cap H_{c,a}$, where $H_0^1(a, +\infty)$ is the usual Sobolev space and

$$H_{c,a} := \left\{ u \in H^{1}_{loc}(a, +\infty) : \int_{a}^{+\infty} e^{ct} u'(t)^{2} dt < +\infty, \quad u(+\infty) = 0 \right\}$$
 (3)

with norm

$$||u||_{c,a} = \int_{a}^{+\infty} e^{ct} \left(u'(t)^2 + u(t)^2 \right) dt$$

and $H_{c,a}^0 := H_{c,a} \cap H_0^1(a, +\infty) = \{u \in H_{c,a} : u(a) = 0\}.$

Equation (1) is obtained by the Fisher-Kolmogorov's equation $u_t = u_{xx} + f(t,u)$, looking for the traveling waves u(x,t) = U(x-ct) with speed c. There is a vast studies on heteroclinic solutions of eq. (1). We refer the reader to Kolmogorov, Petrovsky and Piskunov [3], Aronson and Weinberger [2], Arias and al. [1], Nehari [7], Versini [9], Szulkin [8], Li and Wang [4] and references therain. Solutions of (1) with initial data with compact support are studied in [2]. Fast solutions of Eq. (1) are studied in the paper of Arias and al. [1] via variational methods. Heteroclinic solutions for non-autonomous second order differential equations are studied in [5, 10]. Verzini also studied equation of type (1), when c = 0, and she proves the existence of many oscillating solutions belonging to $L^{\infty}(\mathbb{R})$. She used the variational method and the approach of Nehari [7]. In the present paper we will prove the existence of positive solution of Eq. (1), belonging to $H_0^1(a, +\infty)$, using the methods of [9].

Note that by (H1) it follows

$$\frac{c_1}{q+2}|s|^{q+2} \le F(t,s) \le \frac{c_2}{q+2}|s|^{q+2}, \quad \forall (t,s) \in \mathbb{R}^2, \tag{4}$$

for some positive constants c_1 and c_2 .

We introduce the energy functional associated with the problem (1), (2), further referred as (P):

$$J_{[a,b]}(u) := \int_{a}^{b} e^{ct} \left(\frac{u'(t)^{2}}{2} - F(t, u(t)) \right) dt,$$

$$f^{+\infty} = \int_{a}^{b} u'(t)^{2} dt$$

$$J(u) = J_{[a,+\infty)}(u) := \int_a^{+\infty} e^{ct} \left(\frac{u'(t)^2}{2} - F(t, u(t)) \right) dt,$$

where $a \in \mathbb{R}$ is a fixed real number and $b \in (a, +\infty)$. Let

$$\mu\left(u\right) = \mu_{[a,+\infty)}\left(u\right) := \sup_{\lambda>0} J_{[a,+\infty)}\left(\lambda u\right) = \sup_{\lambda>0} J\left(\lambda u\right). \tag{5}$$

We will prove that for each nonzero function $u, \mu\left(u\right) = \sup_{\lambda>0} J\left(\lambda u\right) \geq C > 0$. There exists unique number $\lambda = \lambda\left(u\right) > 0$, for which $\sup_{\lambda>0} J\left(\lambda u\right)$ is attained, i.e. $\mu\left(u\right) = \sup_{\lambda>0} J\left(\lambda u\right) = J\left(\lambda\left(u\right)u\right)$.

We introduce the following set, analogous to the Nehari manifold:

$$N(a, +\infty) := \{ u \in H_{c,a}^{0} \setminus (0) : \lambda(u) = 1 \}$$
$$= \{ u \in H_{c,a}^{0} \setminus (0) : \nabla J(u) . u = 0 \},$$

as in Verzini [9]. Since we are looking for the solutions of (1), which are non-negative on $[a, +\infty)$, we introduce the set

$$N^{+}(a, +\infty) := \{u \in N (a, +\infty) : u \ge 0\}.$$

Define the function

$$\varphi^{+}\left(a,+\infty\right):=\inf\left\{ \sup_{\lambda>0}J\left(\lambda u\right):u\in H_{c,a}^{0}\smallsetminus\left(0\right),u\geq0\right\} .$$

Our main result is:

Theorem 1. Let the conditions (H1)-(H4) hold.

Then $\varphi^+(a, +\infty)$ is attained by at least one function $u_+ \in N^+(a, +\infty)$, $u_+ > 0$ on $(a, +\infty)$ and $u_+(t)$ is a solution of the problem (P) for $t \in (a, +\infty)$.

The paper is organized as follows. In Section 2 we give preliminaries on the function spaces, embedding inequalities and three lemmas for the corresponding functional J. In Section 3 we give the proof of Theorem 1 and some comments.

2. Preliminaries

By [1] for each $u \in H_{c,a}$ the following inequality holds:

$$\int_{a}^{+\infty} e^{ct} u'(t)^{2} dt \ge \frac{c}{2} e^{ct_{0}} u(t_{0})^{2} + \frac{c^{2}}{4} \int_{a}^{+\infty} e^{ct} u(t)^{2} dt$$
 (6)

for any $t_0 \in [a, +\infty)$. The inequality (6) shows that in the linear space $H_{c,a}$ we can introduce the norm

$$||u||_{H_{c,a}} = \left(\int_{a}^{+\infty} e^{ct} u'(t)^2 dt\right)^{\frac{1}{2}},$$

corresponding to the scalar product $\langle u, v \rangle_{H_{c,a}} = \int_a^{+\infty} e^{ct} u'(t) v'(t) dt$. For each function $u \in H_{c,a} \subset L^{\infty}[a, +\infty)$, $\sup_{t \in [a, +\infty)} |u(t)| < +\infty$ and it is attained, i.e.

$$\sup_{t \in [a, +\infty)} |u(t)| = \max_{t \in [a, +\infty)} |u(t)| = |u(t_1)|$$

for some point $t_1 \in [a, +\infty)$. Further by C we will denote various positive constants not depending on u.

We have the following lemma.

Lemma 1. Let $u \in H_{c,a}$. Then the inequality

$$\int_{a}^{+\infty} e^{ct} u'(t)^{2} dt \ge C(u(t_{1})^{2} + \int_{a}^{+\infty} e^{ct} u(t)^{2} dt + \frac{\int_{a}^{+\infty} e^{ct} |u(t)|^{q+2} dt}{|u(t_{1})|^{q}} + (\int_{a}^{+\infty} e^{ct} |u(t)|^{q+2} dt)^{\frac{2}{q+2}}),$$

$$(7)$$

holds for a constant C > 0.

Proof. We have by (6) that

$$\int_{a}^{+\infty} e^{ct} u'(t)^{2} dt \geq \frac{c}{2} e^{ct_{1}} u(t_{1})^{2},$$

$$\int_{a}^{+\infty} e^{ct} u'(t)^{2} dt \geq \frac{c^{2}}{4} \int_{a}^{+\infty} e^{ct} u(t)^{2} dt.$$

By Young's inequality we get

$$\int_{a}^{+\infty} e^{ct} u'(t)^{2} dt \geq C(u(t_{1})^{2} + \int_{a}^{+\infty} e^{ct} u(t)^{2} dt)$$

$$\geq C \left(u(t_1)^2 + \frac{\int_a^{+\infty} e^{ct} |u(t)|^{q+2} dt}{|u(t_1)|^q} \right)$$

$$\geq C \left(u(t_1)^2 \right)^{\frac{q}{q+2}} \frac{\left(\int_a^{+\infty} e^{ct} |u(t)|^{q+2} dt \right)^{\frac{2}{q+2}}}{\left(|u(t_1)|^q \right)^{\frac{2}{q+2}}}$$

$$= C \left(\int_a^{+\infty} e^{ct} |u(t)|^{q+2} dt \right)^{\frac{2}{q+2}}, \quad C > 0$$

for each nonzero function $u \in H_{c.a}$. These inequalities imply (7).

Next, we have the following

Lemma 2. Let the function $f(t,s) \in C^1(\mathbb{R}^2)$ satisfy the conditions (H1), (H4) and $\mu(u)$ is defined by (5). Then, for every nonzero function $u(t) \in H_{c,a}$, $\mu_{[a,+\infty)}(u) \geq C > 0$, where the constant C does not depend on u.

Proof. By (4),

$$J(\lambda u) = \frac{1}{2}\lambda^{2} \int_{a}^{+\infty} e^{ct}u'(t)^{2} dt - \int_{a}^{+\infty} e^{ct}F(t,\lambda u(t)) dt$$

$$\geq \frac{1}{2}\lambda^{2} \int_{a}^{+\infty} e^{ct}u'(t)^{2} dt - \left(\frac{c_{2}}{q+2} \int_{a}^{+\infty} e^{ct} |u(t)|^{q+2} dt\right) \lambda^{q+2}$$

$$= A_{1}\lambda^{2} - B\lambda^{q+2},$$

where $A_1 = \frac{1}{2} \int_a^{+\infty} e^{ct} u'(t)^2 dt > 0$, $B = \frac{c_2}{q+2} \int_a^{+\infty} e^{ct} |u(t)|^{q+2} dt > 0$, since $u(t) \in H_{c,a}$ is nonzero function. We have that for every $\lambda \in [0, +\infty)$,

$$\mu_{[a,+\infty)}(u) = \sup_{\lambda>0} J(\lambda u)$$

$$\geq c_4(q) \frac{2^{\frac{2}{q}} \cdot q}{(q+2)^{1+\frac{2}{q}}} \left(\frac{\frac{1}{2} \int_a^{+\infty} e^{ct} u'(t)^2 dt}{\left(\frac{c_2}{q+2} \int_a^{+\infty} e^{ct} |u(t)|^{q+2} dt\right)^{\frac{2}{q+2}}} \right)^{\frac{q+2}{q}}$$

$$\geq c_5(q) > 0,$$

where the constants $c_4(q)$ and $c_5(q)$ depend only on q. In the conclusion of the last inequality, we took into account (7). The lemma is proved.

Let $u(t) \in H_{c,a}$ be an arbitrary fixed nonzero function and the conditions of Lemma 2 be fulfilled. By (4) we have

$$J(\lambda u) = \frac{1}{2} \lambda^{2} \int_{a}^{+\infty} e^{ct} u'(t)^{2} dt - \int_{a}^{+\infty} e^{ct} F(t, \lambda u(t)) dt$$

$$\leq \lambda^{2} A_{1} - \left(\frac{c_{1}}{q+2} \int_{a}^{+\infty} e^{ct} |u(t)|^{q+2} dt\right) \lambda^{q+2}$$

$$= A_{1} \lambda^{2} - B_{1} \lambda^{q+2}.$$

Then $J(\lambda u) < 0$ for sufficiently large $\lambda > 0$. Since J(0) = 0 and $J(\lambda u)$ is continuous function in λ , then $\mu(u) = \sup_{\lambda \in J} J(\lambda u)$ is attained.

Moreover $\mu(u) = \mu(ku)$ for every constant k > 0.

Thus for the given nonzero function $u(t) \in H_{c,a}$, there exists a positive number $\lambda_0 > 0$, such that

$$\mu (\lambda_0 u) = J(\lambda_0 u) = \sup_{\lambda > 0} J(\lambda u).$$

If we denote the function $\lambda_0 u \in H_{c,a} \setminus \{0\}$ again by u, then the last equality can be written as

$$\mu\left(u\right) = J\left(u\right). \tag{8}$$

We show that for any nonzero function $v \in H_{c,a}$, i.e. function belonging to $H_{c,a} \setminus \{0\}$, there exists a function $u \in H_{c,a} \setminus \{0\}$ such that u = kv, with suitable constant k > 0, such that (8) holds. As in [9, p.2017] it follows that

$$\frac{\partial}{\partial \lambda} J(\lambda u) \mid_{\lambda=1} = \nabla J(u) u = \int_{a}^{+\infty} e^{ct} \left(u'(t)^{2} - f(t, u(t)) u(t) \right) dt = 0, \quad (9)$$

which holds for critical points of $J(\lambda u)$ as a function of λ .

Lemma 3. Let the function $f(t,s) \in C^1(\mathbb{R}^2)$ satisfy the conditions (H1), (H2) and (H4) and $u \in H_{c,a}$ be nonzero function, for which (9) holds. Then

$$J''(u)[u,u]<0.$$

Moreover there exists unique number $\lambda = \lambda(u) > 0$ such that $\mu(u) = J(\lambda(u)u)$ and the function $u \to \lambda(u)$ is of class C^1 .

Proof. We suppose that for the nonzero function $u \in H_{c,a}$, (9) holds, but $J''(u)[u,u] \geq 0$. Then

$$\int_{a}^{+\infty} e^{ct} \left(u'(t)^{2} - f_{s}(t, u(t)) u(t)^{2} \right) dt \ge 0.$$
 (10)

Subtracting (10) from (9), we obtain that

$$\int_{a}^{+\infty} e^{ct} \left(f_s(t, u(t)) u(t)^2 - f(t, u(t)) u(t) \right) dt \le 0.$$

Taking into account (H1) and (H2), we get

$$0 \geq \int_{a}^{+\infty} e^{ct} \left(f_{s}(t, u(t)) u(t)^{2} - f(t, u(t)) u(t) \right) dt$$

$$\geq \int_{a}^{+\infty} e^{ct} (A - 1) f(t, u(t)) u(t) dt$$

$$\geq c_{5} (A - 1) \int_{a}^{+\infty} e^{ct} |u(t)|^{2+q} dt,$$

where the constant $c_5 > 0$ and A > 1. Thus we proved that $\int_a^{+\infty} e^{ct} |u(t)|^{2+q} dt \le 0$. But it is impossible for nonzero function u(t). The obtained contradiction shows that the considered function u(t) satisfies the inequality J''(u)[u,u] < 0. The rest of the proof of the lemma is as in [9, Proposition 3.1]. Exactly, the unique number $\lambda = \lambda(u) > 0$ such that $\mu(u) = J(\lambda(u)u)$, satisfies the equation

$$\Phi \left(\lambda \left(u\right) ,u\right) :=\nabla J\left(\lambda \left(u\right) u\right) .u=0.$$

Also $\frac{\partial}{\partial \lambda} \Phi\left(\lambda,u\right) = \frac{\partial}{\partial \lambda} \left(\nabla J\left(\lambda u\right).u\right) = J''\left(\lambda u\right)\left[u,u\right] < 0$ for $\lambda = \lambda\left(u\right)$, where Φ is of class C^1 and the function $\lambda = \lambda\left(u\right)$ can be locally implicitly defined. From the implicit function theorem $\lambda\left(u\right)$ is of class C^1 . Lemma 3 is proved.

The considerations in the proof of Lemma 3 show that

$$\varphi^{+}\left(a,+\infty\right) = \inf_{N^{+}\left[a,+\infty\right)} J\left(u\right).$$

3. Proof of the main result

Let $\{u_n\} \subset N^+(a, +\infty)$ be a minimizing sequence for $\varphi^+(a, +\infty)$. Without loss of generality, we can suppose that

$$\varphi^{+}(a, +\infty) + \varepsilon \ge J(u_n) \longrightarrow \varphi^{+}(a, +\infty), \quad n \to +\infty$$
 (11)

for sufficiently small number $\varepsilon > 0$.

Proof of Theorem 1:

Step 1. We have

$$c_6 \int_a^{+\infty} e^{ct} u'_n(t)^2 dt \le J(u_n) \le c_7 \int_a^{+\infty} e^{ct} u'_n(t)^2 dt,$$
 (12)

$$c_8 \int_a^{+\infty} e^{ct} F\left(t, u_n\left(t\right)\right) dt \le J\left(u_n\right) \le c_9 \int_a^{+\infty} e^{ct} F\left(t, u_n\left(t\right)\right) dt \qquad (13)$$

for some positive constants c_i , i = 6, 7, 8, 9. Since $\nabla J(u_n) . u_n = 0$, as in (9),

$$\int_{a}^{+\infty} e^{ct} \left(u_n'(t)^2 - f(t, u_n(t)) u_n(t) \right) dt = 0.$$

We have $J(u_n) = \int_a^{+\infty} e^{ct} \left(\frac{1}{2} u'_n(t)^2 - F(t, u_n(t))\right) dt$. By (H3) and (4),

$$2J(u_n) = \int_a^{+\infty} e^{ct} (f(t, u_n(t)) u_n(t) - 2F(t, u_n(t))) dt$$

$$\geq \alpha \int_a^{+\infty} e^{ct} F(t, u_n(t)) dt \geq \frac{\alpha c_1}{q+2} \int_a^{+\infty} e^{ct} |u_n(t)|^{2+q} dt.$$

This inequality shows that $J(u_n) \geq 0$. Hence

$$\frac{1}{2} \int_{a}^{+\infty} e^{ct} u'_{n}(t)^{2} dt = \int_{a}^{+\infty} e^{ct} F(t, u_{n}(t)) dt + J(u_{n})$$

$$\leq \frac{c_{2}}{q+2} \int_{a}^{+\infty} e^{ct} |u_{n}(t)|^{2+q} dt + J(u_{n})$$

$$\leq \left(1 + \frac{2c_{2}}{\alpha c_{1}}\right) J(u_{n})$$

and

$$J(u_n) \ge \frac{\alpha c_1}{2\alpha c_1 + 4c_2} \int_a^{+\infty} e^{ct} u'_n(t)^2 dt.$$

By the definition of $J(u_n)$ and (4), $J(u_n) \leq \frac{1}{2} \int_a^{+\infty} e^{ct} u_n'(t)^2 dt$. Thus, we proved the first inequality (12) with $c_6 = \frac{\alpha c_1}{2\alpha c_1 + 4c_2}$ and $c_7 = \frac{1}{2}$. The inequality (13) holds with $c_8 = \frac{\alpha}{2}$ and $c_9 = \frac{c_2(q+2)}{2c_1}$ by (H1) and (4) since

$$2J(u_n) \leq \int_a^{+\infty} e^{ct} u'_n(t)^2 dt = \int_a^{+\infty} e^{ct} f(t, u_n(t)) u_n(t) dt$$

$$\leq \frac{c_2(q+2)}{c_1} \int_a^{+\infty} e^{ct} F(t, u_n(t)) dt.$$

These assertions show that $\int_a^{+\infty} e^{ct} u_n'(t)^2 dt$ and $\int_a^{+\infty} e^{ct} |u_n(t)|^{2+q} dt$ are bounded by constant, which does not depend on n. From the inequality (6), the same is true for $\int_a^{+\infty} e^{ct} u_n(t)^2 dt$. Hence

$$\int_{a}^{+\infty} e^{ct} (u_n(t)^2 + u'_n(t)^2) dt$$

and

$$\int_{a}^{+\infty} e^{ct} \left| u_n \left(t \right) \right|^{2+q} dt$$

are bounded by a constant, which does not depend on n. Then, the sequence $\{u_n\}$ is bounded in $H_{c,a}$ equipped by the norm $\|u\|_{H_{c,a}} = \left(\int_a^{+\infty} e^{ct} u'(t)^2 dt\right)^{\frac{1}{2}}$.

Step 2. There exists a function $u_0 \in H_{c,a}^0 := H_{c,a} \cap H_0^1[a, +\infty)$ and subsequence $\{u_{n_k}\}$ still denoted by $\{u_n\}$, such that $u_n \longrightarrow u_0$ for $n \to +\infty$ in the week $H_{c,a}^0$ topology; $u_n \longrightarrow u_0$ for $n \to +\infty$ in the strong L^2 topology and on any bounded and closed subinterval of $[a, +\infty)$. Moreover $u_0 \ge 0$ is nonzero function. Let us remind that $\{u_n\} \subset N^+(a, +\infty)$ and then $u_n \ge 0, n = 1, 2, \ldots$. Hence $u_0 \ge 0$.

We will prove that u_0 is nonzero function. Suppose the contrary, i.e., that $u_0 \equiv 0$. This means that $u_n \longrightarrow 0$ for $n \to +\infty$ in the week $H_{c,a}^0$ topology and $u_n \longrightarrow 0$ for $n \to +\infty$ in the strong L^2 topology on any bounded and closed subinterval of $[a, +\infty)$. Let b > a be an arbitrary number. From [1, p.321],

$$|u_n(b)| \le \left(\frac{e^{-cb}}{c} \int_b^{+\infty} e^{ct} u_n'(t)^2 dt\right)^{\frac{1}{2}} \le c_{10} e^{-\frac{c}{2}b}, \quad \forall b > a$$
 (14)

and the constant $c_{10} > 0$ does not depend on $n \in \mathbb{N}$ and b. Let $b \in (a, +\infty)$ be a fixed (sufficiently large) number. Then for every $\varepsilon > 0$, there exists $n_0 \in \mathbb{N}$, depending on b and ε , such that

$$\int_{a}^{b} e^{ct} u_n(t)^2 dt \le e^{cb} \int_{a}^{b} u_n(t)^2 dt < \varepsilon, \quad \forall n \ge n_0.$$

From (14), replacing b by t, $|u_n(t)| \le c_{10}e^{-\frac{c}{2}t}$, $\forall t \in [a, +\infty)$ and for some $n_1 > n_0$, $n_1 \in \mathbb{N}$, depending on b, we get by q > 0

$$\int_{a}^{b} e^{ct} |u_{n}(t)|^{2+q} dt \le c_{10} \int_{a}^{b} e^{ct} |u_{n}(t)|^{2} dt < \frac{\varepsilon}{2}, \ \forall n \ge n_{1}.$$
 (15)

From (14), we have

$$\int_{b}^{+\infty} e^{ct} |u_n(t)|^{2+q} dt \le \max_{t \in [b, +\infty)} |u_n(t)|^q \cdot \int_{a}^{+\infty} e^{ct} |u_n(t)|^2 dt \le c_{11} e^{-\frac{c}{2}bq} \quad (16)$$

and the constant $c_{11} > 0$ does not depend on $b \in (a, +\infty)$ and $n \in \mathbb{N}$. Now let $\varepsilon > 0$ be a fixed, sufficiently small number. We choose the number $b \in (a, +\infty)$ so large, such that $c_{11}e^{-\frac{c}{2}bq} < \frac{\varepsilon}{2}$. Then by (16),

$$\int_{b}^{+\infty} e^{ct} |u_n(t)|^{2+q} dt < \frac{\varepsilon}{2}, \quad \forall n \in \mathbb{N}.$$

We choose the number $n_1 \in \mathbb{N}$, $n_1 > n_0$, such that (15) holds and hence $\int_a^{+\infty} e^{ct} |u_n(t)|^{2+q} dt < \varepsilon$, $\forall n \geq n_1$. Thus we prove that

$$\lim_{n \to +\infty} \int_{a}^{+\infty} e^{ct} |u_n(t)|^{2+q} dt = 0.$$

By Step 1, (4) and (13) it follows $\lim_{n\to+\infty} J(u_n) = 0$. This contradicts to Lemma 2, according which $J(u_n) = \mu_{[a,+\infty)}(u_n) \geq C > 0$. The contradiction shows that u_0 is a nonzero function.

Step 3. $\int_{a}^{+\infty} e^{ct} u'_{n}(t)^{2} dt \ge \int_{a}^{+\infty} e^{ct} u'_{0}(t)^{2} dt + o(1).$

$$\int_{a}^{+\infty} e^{ct} u'_{n}(t)^{2} dt = \int_{a}^{+\infty} e^{ct} \left[u'_{0}(t) + \left(u'_{n}(t) - u'_{0}(t) \right) \right]^{2} dt \qquad (17)$$

$$= \int_{a}^{+\infty} e^{ct} u'_{0}(t)^{2} dt + 2 \int_{a}^{+\infty} e^{ct} \left(u'_{n}(t) - u'_{0}(t) \right) u'_{0}(t) dt + \int_{a}^{+\infty} e^{ct} \left(u'_{n}(t) - u'_{0}(t) \right)^{2} dt.$$

By $u_n \longrightarrow u_0$ for $n \to +\infty$ weakly in $H_{c,a}^0$ we have

$$\lim_{n \to +\infty} \int_{a}^{+\infty} e^{ct} \left(u'_{n}(t) - u'_{0}(t) \right) u'_{0}(t) dt$$

$$= \lim_{n \to +\infty} \left\langle u'_{n}(t) - u'_{0}(t), u'_{0}(t) \right\rangle_{H_{c,a}} = 0.$$

Then

$$2\int_{a}^{+\infty}e^{ct}\left(u_{n}^{\prime}\left(t\right)-u_{0}^{\prime}\left(t\right)\right)u_{0}^{\prime}\left(t\right)dt=o\left(1\right),$$

and by (17),

$$\int_{a}^{+\infty} e^{ct} u'_{n}(t)^{2} dt \ge \int_{a}^{+\infty} e^{ct} u'_{0}(t)^{2} dt + o(1).$$

Step 4. $\lim_{n\to+\infty} \int_a^{+\infty} e^{ct} F(t, u_n(t)) dt = \int_a^{+\infty} e^{ct} F(t, u_0(t)) dt$.

As in Step 2, replacing u_n by $u_n - u_0$, we can prove that

$$\lim_{n \to +\infty} \int_{a}^{+\infty} e^{ct} |u_n(t) - u_0(t)|^{2+q} dt = 0.$$
 (18)

By (H1) and
$$u_n(t) \ge 0$$
, $u_0(t) \ge 0$,

$$|F(t, u_n(t)) - F(t, u_0(t))| = \left| \int_{u_0(t)}^{u_n(t)} f(t, \tau) d\tau \right| \le c_2 \left| \int_{u_0(t)}^{u_n(t)} \tau^{1+q} d\tau \right|$$
$$= \frac{c_2}{q+2} \left| u_n(t)^{2+q} - u_0(t)^{2+q} \right|.$$

Hence

$$\left| \int_{a}^{+\infty} e^{ct} F(t, u_{n}(t)) dt - \int_{a}^{+\infty} e^{ct} F(t, u_{0}(t)) dt \right|$$

$$\leq \frac{c_{2}}{q+2} \int_{a}^{+\infty} e^{ct} \left| u_{n}(t)^{2+q} - u_{0}(t)^{2+q} \right| dt$$

$$\leq K_{1} \int_{a}^{+\infty} e^{ct} \left| u_{n}(t) - u_{0}(t) \right| \left(u_{n}(t)^{1+q} + u_{0}(t)^{1+q} \right) dt.$$

Using the Hölder inequality, it is easy to obtain that

$$\left| \int_{a}^{+\infty} e^{ct} F(t, u_n(t)) dt - \int_{a}^{+\infty} e^{ct} F(t, u_0(t)) dt \right|$$

$$\leq K_1 \left(\int_{a}^{+\infty} e^{ct} (u_n(t)^{1+q} + u_0(t)^{1+q})^{\frac{2+q}{1+q}} dt \right)^{\frac{1+q}{2+q}}$$

$$\left(\int_{a}^{+\infty} e^{ct} |u_n(t) - u_0(t)|^{2+q} dt \right)^{\frac{1}{2+q}}$$

$$\leq K_{2} \left(\int_{a}^{+\infty} e^{ct} u_{n}(t)^{2+q} dt + \int_{a}^{+\infty} e^{ct} u_{0}(t)^{2+q} dt \right)^{\frac{1+q}{2+q}}$$

$$\left(\int_{a}^{+\infty} e^{ct} |u_{n}(t) - u_{0}(t)|^{2+q} dt \right)^{\frac{1}{2+q}}$$

$$\leq K_{3} \left(\int_{a}^{+\infty} e^{ct} |u_{n}(t) - u_{0}(t)|^{2+q} dt \right)^{\frac{1}{2+q}},$$

where K_j , j = 1, 2, 3 are constants not depending on u and n. Taking into account (18), we obtain the assertion of Step 4.

Step 5. There exists a nonzero function $u_{+} \in N^{+}(a, +\infty)$, which $u_{+} \geq 0$ and $J_{[a,+\infty)}(u_{+}) = \varphi^{+}(a,+\infty) = \inf_{N^{+}(a,+\infty)} J_{[a,+\infty)}(u) > 0$.

From (4), Step 3 and Step 4 it follows that $J(u_n) \geq J(u_0) + o(1)$ (i.e., the functional J is weakly lower semi continuous). Hence the inequality $J_{[a,+\infty)}(\lambda u_0) \leq J_{[a,+\infty)}(\lambda u_n) + o(1)$, $\forall \lambda > 0$ holds.

Define

$$u_+ := \lambda (u_0) u_0.$$

We obtain, that u_+ satisfies the conditions of Step 5.

We will prove that $u_+ > 0$ on $(a, +\infty)$. For this purpose, we adapt the proof of Theorem 3.1 of [9, pp. 2019-2020] to our case.

Claim 1. $u_+ \in C^1(a, +\infty)$. We assume that there exists $\tau \in (a, +\infty)$, for which $u_+(\tau) = 0$ and $u'_+(t)$ is not continuous for $t = \tau$. Let $u'_+(\tau - 0) < 0$ and the constants $\rho > 0$ and $\varepsilon > 0$ are sufficiently small numbers. For $\lambda \in [1 - \varepsilon, 1 + \varepsilon]$ we consider the class of problems

$$\inf \left\{ J_{\left[\tau-\rho,\tau+\rho\right]}(\upsilon) : \upsilon \in H^{1}\left(\tau-\rho,\tau+\rho\right), \quad \upsilon\left(\tau+\rho\right) = \lambda u_{+}\left(\tau-\rho\right) \\ \left\|\upsilon\right\|_{\infty} \leq 1 \right\}$$
 (19)

where $\tau - \rho, \tau + \rho \in (a, +\infty)$. We have

$$\frac{d^{2}}{d\lambda^{2}}J_{[\tau-\rho,\tau+\rho]}(u+\lambda\varphi)_{\lambda=0} = \int_{\tau-\rho}^{\tau+\rho} e^{ct} \left(\varphi'(t)^{2} - f'_{s}(t,u(t))\varphi(t)^{2}\right) dt \qquad (20)$$

$$\geq \int_{\tau-\rho}^{\tau+\rho} e^{ct}\varphi'(t)^{2} dt - c_{13} \int_{\tau-\rho}^{\tau+\rho} e^{ct}\varphi(t)^{2} dt$$

$$\geq \left(\frac{c_{14}}{\rho^{2}} - c_{3}\right) \int_{\tau-\rho}^{\tau+\rho} e^{ct}\varphi(t)^{2} dt > 0,$$

where the constants $c_{13} := \sup \{f'_s(t,u) : a \leq t < +\infty, -1 \leq u \leq 1\}, c_{14} > 0$, and the function $\varphi(t)$ vanishes on at least one point $t = t_0 \in [\tau - \rho, \tau + \rho]$. As in [9, p. 2026, Lemma 5.1], we can conclude, that J is strictly convex, and thus the minimum of (19) is uniquely achieved by a function v_{λ} . We will show that v_{λ} satisfies Eq. (1.) and $\lim_{\rho \to 0} \|v_{\lambda}\|_{H^1(\tau - \rho, \tau + \rho)} = 0$. For this purpose, we need to prove that $|v_{\lambda}(t)| < 1$ and $v_{\lambda}(t) \geq 0$ for every $t \in [\tau - \rho, \tau + \rho]$.

First, we prove that $\|v_{\lambda}\|_{\infty} = \max_{t \in [\tau - \rho, \tau + \rho]} |v_{\lambda}(t)| < 1$. Suppose the contrary, $\|v_{\lambda}\|_{\infty} = 1$. By inclusion $H^{1}(\tau - \rho, \tau + \rho) \subset C[\tau - \rho, \tau + \rho]$, $u_{+} \in C[\tau - \rho, \tau + \rho]$ and $v_{\lambda} \in C[\tau - \rho, \tau + \rho]$. By $u_{+}(\tau) = 0$, $v_{\lambda}(\tau \pm \rho) = \lambda u_{+}(\tau \pm \rho) = o(1)$ for $\rho \to 0+$. If $v_{\lambda}(\tau_{1}) = \pm 1$ for some $\tau_{1} \in (\tau - \rho, \tau + \rho)$, then

$$|v_{\lambda}(\tau_{1}) - v_{\lambda}(\tau - \rho)| = 1 - o(1) = \left| \int_{\tau - \rho}^{\tau_{1}} v_{\lambda}'(t) dt \right|$$

$$\leq \int_{\tau - \rho}^{\tau + \rho} e^{-\frac{ct}{2}} e^{\frac{ct}{2}} |v_{\lambda}'(t)| dt$$

$$\leq \left(\int_{\tau - \rho}^{\tau + \rho} e^{-ct} dt \right)^{\frac{1}{2}} \left(\int_{\tau - \rho}^{\tau + \rho} e^{ct} v_{\lambda}'(t)^{2} dt \right)^{\frac{1}{2}}$$

$$= \left(\frac{e^{-c(\tau - \rho)} - e^{-c(\tau + \rho)}}{c} \right)^{\frac{1}{2}} \left(\int_{\tau - \rho}^{\tau + \rho} e^{ct} v_{\lambda}'(t)^{2} dt \right)^{\frac{1}{2}}$$

$$\leq c_{15} \sqrt{\rho} \left(\int_{\tau - \rho}^{\tau + \rho} e^{ct} v_{\lambda}'(t)^{2} dt \right)^{\frac{1}{2}},$$

where the constant $c_{15} > 0$ is close to $(2e^{-c\tau})^{\frac{1}{2}}$ for small ρ , i.e. c_{15} depends only on τ . Since $1 - o(1) \ge \frac{1}{\sqrt{2}}$ for small ρ , then

$$\int_{\tau-\rho}^{\tau+\rho} e^{ct} v_{\lambda}'(t)^2 dt \ge \frac{1}{2c_{15}^2 \rho}.$$

Since $F(t, v_{\lambda})$ is bounded for $t \in [\tau - \rho, \tau + \rho]$, and $\frac{1}{2c_{15}^2\rho} - C \ge \frac{c_{16}}{\rho}$ for small $\rho > 0$, it implies

$$J_{[\tau-\rho,\tau+\rho]}(v_{\lambda}) \ge \frac{c_{16}}{\rho}.$$
(21)

Remind that $u_{+} \in H^{1}(\tau - \rho, \tau + \rho)$ and $u_{+}(\tau) = 0$. As above

$$\max_{[\tau-\rho,\tau+\rho]} |u_{+}(t)| = \max_{[\tau-\rho,\tau+\rho]} |u_{+}(t) - u_{+}(\tau)|$$

$$\leq c_{15}\sqrt{\rho} \left(\int_{\tau-\rho}^{\tau+\rho} e^{ct} u'_{+}(t)^{2} dt \right)^{\frac{1}{2}} = o(1)\sqrt{\rho},$$

because $u_+ \in H^0_{c,a}$ and then $\left(\int_{\tau-\rho}^{\tau+\rho} e^{ct} u'_+(t)^2 dt\right)^{\frac{1}{2}} = o(1)$ when $\rho > 0$ is sufficiently small. Thus

$$|\lambda u_{+}(\tau \pm \rho)| = o(\sqrt{\rho}), \ \forall \lambda \in [1 - \varepsilon, 1 + \varepsilon].$$

Now we consider the linear function

$$w(t) = u_{+}(\tau - \rho) - \frac{u_{+}(\tau - \rho) - u_{+}(\tau + \rho)}{2\rho}(t - \tau + \rho).$$

Evidently

$$w(t-\rho) = u_{+}(\tau - \rho),$$

 $w\left(t+\rho\right)=u_{+}\left(\tau+\rho\right)$ and $w'\left(t\right)=\frac{u_{+}\left(\tau+\rho\right)-u_{+}\left(\tau-\rho\right)}{2\rho}.$ As above we have

$$|u_{+}(\tau - \rho) - u_{+}(\tau + \rho)| = \left| \int_{\tau - \rho}^{\tau + \rho} u'(s) \, ds \right|$$

$$\leq c_{15} \sqrt{\rho} \left(\int_{\tau - \rho}^{\tau + \rho} e^{ct} u'_{+}(t)^{2} \, dt \right)^{\frac{1}{2}} = \sqrt{\rho}.o(1),$$

because $u_+ \in H^0_{c,a}$. Then $\left(\int_{\tau-\rho}^{\tau+\rho} e^{ct} u'_+(t)^2 dt\right)^{\frac{1}{2}} = o(1)$ when $\rho > 0$ is sufficiently small. Thus

$$\begin{aligned} \left|w'\left(t\right)\right| &\leq & \frac{C}{\sqrt{\rho}}.o\left(1\right), \\ 0 &\leq & \int_{\tau-\rho}^{\tau+\rho} e^{ct} w'\left(t\right)^{2} dt \leq \frac{C}{\rho}.o\left(1\right) \int_{\tau-\rho}^{\tau+\rho} e^{ct} dt \leq o\left(1\right). \end{aligned}$$

Also,

$$|w(t)| \leq \max (u_{+}(\tau - \rho), u_{+}(\tau + \rho))$$

$$\leq \max_{[\tau - \rho, \tau + \rho]} |u_{+}(t)| = o(1)\sqrt{\rho}, \quad \forall t \in [\tau - \rho, \tau + \rho].$$

Hence $\|w\|_{H^{1}[\tau-\rho,\tau+\rho]} = o(1)$ and

$$J_{\left[\tau-\rho,\tau+\rho\right]}\left(w\right)=o\left(1\right),\ J_{\left[\tau-\rho,\tau+\rho\right]}\left(\lambda w\right)=o\left(1\right)\ \forall\lambda\in\left[1-\varepsilon,1+\varepsilon\right].\tag{22}$$

From (21) and (22), $J_{[\tau-\rho,\tau+\rho]}(\lambda w) \ll J_{[\tau-\rho,\tau+\rho]}(v_{\lambda})$, $\forall \lambda \in [1-\varepsilon, 1+\varepsilon]$ for sufficiently small $\rho > 0$, (where $a \ll b$ means that $\frac{a}{b} = o(1)$). The last inequality contradicts to the fact, that infimum in (19) is attained by the function v_{λ} . The contradiction is due to the assumption, that $||v_{\lambda}||_{\infty} = 1$. Therefore $||v_{\lambda}||_{\infty} \ll 1$ and thus v_{λ} is a solution of Eq.(1).

Now we will prove that $v_{\lambda} \geq 0$ for $t \in [\tau - \rho, \tau + \rho]$. Suppose the contrary. Then changing the sign, v_{λ} vanish at some point of $[\tau - \rho, \tau + \rho]$ and by $v_{\lambda}(\tau \pm \rho) = \lambda u_{+}(\tau \pm \rho) \geq 0$, it follows that v'_{λ} also vanish at some other point of $[\tau - \rho, \tau + \rho]$. Since v_{λ} satisfies (1), then

$$\left(e^{ct}v_{\lambda}'\right)' + e^{ct}f\left(t, v_{\lambda}\right) = 0 \Longrightarrow v_{\lambda}'\left(t\right) = -\int_{t_{1}}^{t} e^{-c(t-s)}f\left(s, v_{\lambda}\left(s\right)\right)ds,$$

where $t_1 \in [\tau - \rho, \tau + \rho]$ is such that $\upsilon'_{\lambda}(t_1) = 0$. Then

$$\begin{aligned} \left| v_{\lambda}'(t) \right| & \leq \int_{t_{1}}^{t} e^{-c(t-s)} \left| f\left(s, v_{\lambda}\left(s \right) \right) \right| ds \\ & \leq c_{17} \int_{\tau-\rho}^{\tau+\rho} e^{-c(t-s)} ds = c_{17} \frac{e^{-c(\tau-\rho)} - e^{-c(\tau+\rho)}}{c} \leq c_{18} \rho, \end{aligned}$$

because $f\left(s, v_{\lambda}\left(s\right)\right)$ is bounded, for $s \in \left[\tau - \rho, \tau + \rho\right]$. Thus

$$\left| v_{\lambda}'(t) \right| \le c_{18}\rho \ \forall t \in \left[\tau - \rho, \tau + \rho\right],\tag{23}$$

where the constant $c_{18} > 0$ depends only on τ . Recall that $u_+(\tau) = 0$ and $u'_+(\tau - 0) < 0$. Then for $\rho > 0$ sufficiently small, $u_+(\tau - \rho) = u_+(\tau) - \rho u'_+(\tau - 0) = -\rho u'_+(\tau - 0)$. Hence

$$\upsilon_{\lambda} (\tau - \rho) = \lambda u_{+} (\tau - \rho) = -\lambda \rho u'_{+} (\tau - 0),$$

i.e. $c_{19}\rho \leq v_{\lambda}(\tau - \rho) \leq c_{20}\rho$ for $\rho > 0$ sufficiently small. Then $v_{\lambda}(t) = v_{\lambda}(\tau - \rho) + \int_{\tau - \rho}^{t} v_{\lambda}'(s) ds$ and from (23)

$$\left| \int_{\tau-\rho}^{t} \upsilon_{\lambda}'(s) \, ds \right| \le c_{18} \rho^{2} \ll \upsilon_{\lambda} \left(\tau-\rho\right), \quad \forall t \in \left[\tau-\rho, \tau+\rho\right].$$

Thus $v_{\lambda}(t)$ cannot change the sign when $t \in [\tau - \rho, \tau + \rho]$. This consideration is true also, when $u'_{+}(\tau - 0) = -\infty$, because $v_{\lambda}(\tau - \rho) \geq C\rho$ for sufficiently large constant C > 0 and $v_{\lambda}(t)$ does not change the sign again. We prove, that $v_{\lambda} \geq 0$ for $t \in [\tau - \rho, \tau + \rho]$.

Note that $v_{\lambda}(t) \in C^{1}(\tau - \rho, \tau + \rho)$, because $v_{\lambda}(t)$ satisfies Eq. (1) for $t \in (\tau - \rho, \tau + \rho)$. Now we define the function

$$\widetilde{u}_{\lambda}\left(t\right):=\left\{\begin{array}{ll}v_{\lambda}\left(t\right), & t\in\left[\tau-\rho,\tau+\rho\right],\\ \lambda u_{+}\left(t\right), & t\in\left[a,+\infty\right)\smallsetminus\left[\tau-\rho,\tau+\rho\right].\end{array}\right.$$

We should note that $\widetilde{u}_{\lambda}(t) \geq 0 \quad \forall t \in [a, +\infty)$ and $\widetilde{u}_{\lambda}(t) \in H_{c,a}^{0}$, because $v_{\lambda}(t)$ satisfies the boundary conditions in (19). We will prove, that $\widetilde{u}_{\lambda} \to \lambda u_{+}$ in $H_{c,a}^{0}$, when $\rho \to 0$. Indeed, from (22), $J_{[\tau-\rho,\tau+\rho]}(v_{\lambda}) \leq J_{[\tau-\rho,\tau+\rho]}(w) \leq o(1)$ and since

$$\|v_{\lambda}\|_{\infty} < 1 \Longrightarrow \int_{\tau-\rho}^{\tau+\rho} F\left(t, v_{\lambda}\left(t\right)\right) dt < \frac{2c_{2}}{q+2}\rho,$$

then $\int_{\tau-\rho}^{\tau+\rho} e^{ct} v_{\lambda}'(t)^2 dt = o(1)$ when $\rho \to 0$. This proves that $u_{\lambda} \to \lambda u_+$ in $H_{c,a}^0$, when $\rho \to 0$. Further the proof that $u_+ \in C^1(a, +\infty)$ holds as in [9, p.2020]. Finally we should note that there exists finite right derivative $u'_+(a+0)$. Indeed,

if $u_{+}(t) \equiv 0$ in a small right neighborhood of a, then obviously $u'_{+}(a+0) = 0$. If $u_{+}(t) > 0$ in a small right neighborhood of a, then $u_{+}(t)$ satisfies (1) and it is easy to show that

$$u'_{+}(a+0) := \lim_{t \to a+0} u'_{+}(t)$$

$$= e^{-c(a-t_{0})} \left[u'_{+}(t_{0}) + \int_{a}^{t_{0}} e^{c(s-t_{0})} f(s, u_{+}(s)) ds \right],$$

where $t_0 > a$ is a point, close enough to a. Thus we prove, that $u_+ \in C^1[a, +\infty)$.

Claim 2. $u_{+}(t) > 0, \forall t \in [a, +\infty)$.

We assume that $u_+(t_0) = 0$ for some $t_0 > a$, and $u_+(t) > 0$ in a small right (left) neibourghood of t_0 . Since $u_+ \ge 0$, then t_0 is a point of local minimum for u_+ and since $u_+ \in C^1(a, +\infty)$, then $u'_+(t_0) = 0$. Since $f(t, 0) \equiv 0$, then $u \equiv 0$ is a solution of (1) for the Cauchy conditions $u(t_0) = u'(t_0) = 0$ and from the uniqueness theorem of the local Cauchy problem, follows that $u_+(t) \equiv 0$ in a small neibourhood of t_0 . Thus it implies $u_+(t) = 0$, $\forall t \in [a, +\infty)$, which contradicts to the definition of u_+ as a nonzero function. This contradiction proves the assertion of Claim 2. Theorem 1 is proved.

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